## **Chalmers presentations**

- Per-Simon Kildal (40 min)
  - (Tutorial) Characterization of Feeds for Radio Telescopes (including combination feeds)
- 12:30 Lunch
- Per-Simon Kildal (15min)
  - Introduction to Design of Eleven Feed for SKA and VLBI2010
  - Selection of Geometry and Port lay-out
- Leif Helldner (15min)
  - Cryogenic and Mechanical Design for the Eleven Feed
- Jian Yang (15min)
  - Numerical Optimization of log-Period Antennas and Measurements
- Per-Simon Kildal (10min)
  - Summary of Performance and What Next

Characterization of Feeds for Radio Telescopes

Per-Simon Kildal Chalmers University of Technology 41296 Gothenburg, Sweden www.kildal.se Traditional characterization not good enough for diagnosis

- Traditionally feeds where characterized only by the feed illumination taper in Eand H-planes, and the E- and H-plane phase centers
- This is OK for traditional rotationally symmetric feeds, but NOT for modern compact or wideband feeds
- Then, we need to a better diagnosis approach

## **Content of tutorial**

- Radiation field function (phase ref point)
- BOR antennas and BOR1 efficiency
- Subtended half angle versus F/D
- Feed efficiency and subefficiencies
  - Polarization, Spillover, Illumination, Phase (phase center)
- Blockage and subreflector diffraction
- More efficiencies
  - Mismatch factor, radiation efficiency
- How do they affect G/Tsys
- Combination feeds and decoupling efficiency

# Radiation field and radiation field function



When 
$$r = \sqrt{x^2 + y^2 + z^2}$$
, the phase reference point of  $G(\theta, \phi)$  is origin (0,0,0)

### Spherical coordinate system



 $\mathbf{r} = r\sin\theta\cos\varphi \hat{\mathbf{x}} + r\sin\theta\sin\varphi \hat{\mathbf{y}} + r\cos\theta \hat{\mathbf{z}}$ 

# Equivalent forms of radiation field and its function

$$\mathbf{E}(\mathbf{r}) = \frac{1}{r} e^{-jkr} \mathbf{G}(\hat{\mathbf{r}}) \qquad \mathbf{E}(r, \theta, \varphi) = \frac{1}{r} e^{-jkr} \mathbf{G}'(\theta, \varphi)$$

$$\mathbf{r} = \mathbf{x}\hat{\mathbf{x}} + \mathbf{y}\hat{\mathbf{y}} + \mathbf{z}\hat{\mathbf{z}} = \hat{\mathbf{r}}\mathbf{r}$$

 $\mathbf{r} = r\sin\theta\cos\varphi\,\hat{\mathbf{x}} + r\sin\theta\sin\varphi\,\hat{\mathbf{y}} + r\cos\theta\,\hat{\mathbf{z}}$ 

$$G'(\theta, \varphi) = G(\hat{\mathbf{r}})$$
  
direction  $\hat{\mathbf{r}}$ 

#### Polarization and polarization vectors

Co- and cross-polar radiation field functions:

$$G_{co}(\theta, \varphi) = \boldsymbol{G}(\theta, \varphi) \cdot \hat{\boldsymbol{co}}^{*}(\theta, \varphi)$$
$$G_{xp}(\theta, \varphi) = \boldsymbol{G}(\theta, \varphi) \cdot \hat{\boldsymbol{xp}}^{*}(\theta, \varphi)$$

Convenient choice for<br/>linear y-polarization, i.e.<br/>Ludwig's third definition: $\hat{co}(\theta, \phi) = \hat{y}'(\theta, \phi) = \sin\phi\hat{\theta} + \cos\phi\hat{\phi}$  $\hat{xp}(\theta, \phi) = \hat{x}'(\theta, \phi) = \cos\phi\hat{\theta} - \sin\phi\hat{\phi}$ 

For RHC pol.:

$$\hat{\boldsymbol{co}}(\theta, \varphi) = [\hat{\boldsymbol{x}}'(\theta, \varphi) - j\hat{\boldsymbol{y}}'(\theta, \varphi)]/\sqrt{2} = e^{-j\phi}[\hat{\theta} - j\hat{\varphi}]/\sqrt{2}$$
$$\hat{\boldsymbol{xp}}(\theta, \varphi) = [\hat{\boldsymbol{x}}'(\theta, \varphi) + j\hat{\boldsymbol{y}}'(\theta, \varphi)]/\sqrt{2} = e^{j\phi}[\hat{\theta} + j\hat{\varphi}]/\sqrt{2}$$

# Fourier expansion of radiation field function

• The azimuth variation of the pattern can be expanded in a Fourier series in two ways, :

$$\hat{G}(\theta, \varphi) = G_{co}(\theta, \varphi)\hat{co} + G_{xp}(\theta, \varphi)\hat{xp} =$$

$$= \sum_{n=1}^{\infty} [CO_{sn}(\theta) \sin(n\varphi) + CO_{cn}(\theta) \cos(n\varphi)]\hat{co} + \sum_{n=1}^{\infty} [XP_{sn}(\theta) \sin(n\varphi) + XP_{cn}(\theta) \cos(n\varphi)]\hat{xp}$$

• Alternative which is consistent with spherical TE and TM mode expansions (note that  $A_n(0) = B_n(0) = C_n(0) = D_n(0) = 0$  for  $\theta = 0$  except when n = 1

$$G(\theta, \phi) = G_{\theta}(\theta, \phi)\hat{\theta} + G_{\phi}(\theta, \phi)\hat{\phi} =$$

$$\sum_{n=1}^{\infty} [A_{n}(\theta)\sin(n\phi) + B_{n}(\theta)\cos(n\phi)]\hat{\theta} + \sum_{n=1}^{\infty} [C_{n}(\theta)\cos(n\phi) - D_{n}(\theta)\sin(n\phi)]\hat{\phi}$$

## **Power integral**

• Total radiated power:

$$\mathsf{P}_{\mathsf{rad}} = \frac{1}{2\eta} \int_{4\pi} [|\mathsf{G}_{\varphi}(\theta, \varphi)| + |\mathsf{G}_{\theta}(\theta, \varphi)|] \sin\theta d\theta d\varphi$$

• We will define a power integral:  $P = \iint_{4\pi} [|G_{\phi}(\theta, \phi)|^{2} + |G_{\theta}(\theta, \phi)|^{2}] \sin\theta d\theta d\phi$ 

- The power integral can also be calculated by:
  - aperture integration (for some aperture antennas)
  - integration over feed pattern (in reflector antennas)

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### **Directivity and BOR1 efficiency**

- Power integral:  $P = \sum_{n \in \mathbb{N}} P_n$
- Directivity  $D = \frac{4\pi \left| G_{co}(0) \right|^2}{P} = e_{BOR1} D_{BOR1}$
- BOR1 efficiency

$$e_{BOR1} = \frac{P_1}{P}$$

• BOR1 directivity

$$D_{BOR1} = \frac{4\pi \left| G_{co}(0) \right|^2}{P_1}$$

### **BOR** antennas

- We use the term BOR (Bodies Of Revolution) antennas to characterize antennas which material parts are completely rotationally symmetric around the same (z-)axis.
- The field can still have an azimuth variation. The order of this variation is characterized by an index on BOR.
- The abbreviation BOR is well known in MoM analysis.





Reflector antenna with dipole-disk feed

Corrugated horn antenna

# BOR<sub>0</sub> antennas

- BOR = Body of revolution
- BOR<sub>0</sub> antennas (no azimuth variation):
- Electric dipole excited:
- Magnetic dipole excited:  $G_{e}(\theta, \phi) B_{0}(\theta)\hat{\theta}$

 $G_{e}(\theta, \phi) - B_{0}(\theta)\hat{\theta}$  $G_{m}(\theta, \phi) = C_{o}(\theta)\hat{\phi}$ 

• Example:



Biconical BOR<sub>0</sub> antenna

### BOR<sub>1</sub> antennas (excited for single phi-variation only)

BOR with y-polarized TE<sub>11</sub> type of excitation:

 $G(\theta, \phi) = G_{E}(\theta) \sin \phi \hat{\theta} + G_{H}(\theta) \cos \phi \hat{\phi}$   $\uparrow_{E-plane pattern} \uparrow_{H-plane pattern}$ 

Co- and crosspolar components:

$$G_{co}(\theta, \phi) = G_{y}(\theta, \phi) \cdot \hat{co}^{*} = G_{co45}(\theta) - G_{xp45}(\theta) \cos 2\phi$$
  
$$G_{xp}(\theta, \phi) = G_{y}(\theta, \phi) \cdot \hat{xp}^{*} = G_{xp45}(\theta) \sin 2\phi$$

Co- and crosspolar components in 45 deg plane:

 $G_{co45}(\theta) = \frac{1}{2}[G_{E}(\theta) + G_{H}(\theta)]$  $G_{xp45}(\theta) = \frac{1}{2}[G_{E}(\theta) - G_{H}(\theta)]$ 

The BOR1 relations

We can construct the whole radiation pattern if we know the coand crosspolar patterns in the 45 deg plane, or the E- and Hplane patterns.

### BOR<sub>1</sub> antenna: Circular pol. (RHC)

• RHC co- and cross-polar patterns become:

$$G_{co}(\theta, \varphi) = G_{c} \cdot \hat{co^{*}} = G_{co45}(\theta)$$
$$G_{xp}(\theta, \varphi) = G_{c} \cdot \hat{xp^{*}} = G_{xp45}(\theta)e^{-j2\varphi}$$

• Where as before:

$$G_{co45}(\theta) = \frac{1}{2}[G_{E}(\theta) + G_{H}(\theta)] \qquad G_{xp45}(\theta) = \frac{1}{2}[G_{E}(\theta) - G_{H}(\theta)]$$

The co- and xp- 45 deg patterns for linear polarization are equal to the co- and xp-patterns in all planes for circular polarization. Useful relation, but valid only for ideal BOR1 antennas.

#### SOFT choke horn feed: Constant beamwidth over 0.9-1.7 GHz (Ying, Kishk, Kildal, 1995)



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### Cassegrain antenna

Paraboloidal main reflector and hyperbloidal subreflector. Four independent parameters  $D, d, \theta_0, \psi_0$ Primary focal point, secondary focal point.

![](_page_18_Figure_2.jpeg)

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# Paraboloid: Aperture or feed efficiency

$$e_{ap} = \frac{4\pi \cot^{2}(\theta_{0}/2) \left| \int_{0}^{\theta_{0}} G_{c045}(\theta_{f}) \tan(\theta_{f}/2) d\theta_{f} \right|^{2}}{2\pi \int_{0}^{\pi} [|G_{c045}(\theta_{f})|^{2} + |G_{xp45}(\theta_{f})|^{2}] \sin\theta_{f} d\theta_{f}}$$

Alternative expression based on aperture integration:

$$e_{ap} = e_{sp} \left( \frac{\left(\frac{1}{A}\right) \left| 2\pi \int_{0}^{D/2} E_{co \, 45}(\rho) \rho \, d\rho \right|^{2}}{\frac{D/2}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} + \left| E_{xp \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} \right] \rho \, d\rho \frac{1}{2\pi \int_{0}^{D/2} \left[ \left| E_{co \, 45}(\rho) \right|^{2} \right] \rho$$

Similar expressions exist for reflector antennas of other forms.

### Paraboloid: Spillover efficiency

![](_page_21_Figure_1.jpeg)

### Cassegrain antenna

Paraboloidal main reflector and hyperbloidal subreflector. Four independent parameters, d,  $\theta_0$ ,  $\psi_0$ Primary focal point, secondary focal point.

![](_page_22_Figure_2.jpeg)

# Subefficiencies of Paraboloids and Cassegrain Antennas

Similar formulas apply to general multi-reflector systems.

Factorization of feed efficiency:  $e_{ap} = e_{sp}e_{pol}e_{ill}e_{\phi}$ 

Spillover, polarization, illumination and phase eff.

Spillover efficiency  $e_{sp}$ Relative spillover power is given by  $1 - e_{sp}$ Typically between -0.05 dB and -0.5 dB. Major contributor to the antenna noise temperature.

![](_page_24_Figure_0.jpeg)

![](_page_25_Figure_0.jpeg)

## Paraboloid and Cassegrain: Phase efficiency

$$e_{\phi} = \frac{\left| \int_{0}^{D/2} E_{co45}(\rho) \rho d\rho \right|^{2}}{\left[ \int_{0}^{D/2} |E_{co45}(\rho)| \rho d\rho \right]^{2}} = \frac{\left| \int_{0}^{\theta_{0}} G_{co45}(\theta_{f}) \tan(\theta_{f}/2) d\theta_{f} \right|^{2}}{\left[ \int_{0}^{\theta_{0}} |E_{co45}(\rho)| \rho d\rho \right]^{2}}$$

Typically better than 0.1 dB.

Use to define a <u>phase center</u> for the feed As the phase reference point that maximizes phas efficiency.

Paraboloid and Cassegrain: Maximum displacemen t of phase center from focal point (for -0.1 dB phase eff.) versus subtended half angle

![](_page_27_Figure_1.jpeg)

#### Hat feed in ring-focus paraboloid 1986

![](_page_28_Picture_1.jpeg)

![](_page_28_Picture_2.jpeg)

#### Low sidelobes Good efficiency

# 2007: Master student Denstedt improved bandwidth of hat feed from 10 to 32%

![](_page_29_Picture_1.jpeg)

![](_page_29_Figure_2.jpeg)

Inherently this type of antenna has very low polarization and phase efficiency. Reflector shaped for high phase efficiency. Hat tuned for both good impedance match and high polarization efficiency.

![](_page_29_Picture_4.jpeg)

![](_page_30_Figure_0.jpeg)

#### Improved sector-shaped radiation pattern of horn with shaped lens

![](_page_31_Picture_1.jpeg)

![](_page_31_Picture_2.jpeg)

![](_page_31_Figure_3.jpeg)

![](_page_31_Figure_4.jpeg)

# The horn with shpaed lens should give 0.5 dB higher illumunation efficiency than a horn with a normal feed pattern, but it did not.

![](_page_32_Figure_1.jpeg)

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#### Geometrical aperture blockage model is often used, but does not work in small reflectors, see next slide

![](_page_34_Figure_1.jpeg)

### Blockage efficiency

The total efficiency with blockage can be written:

![](_page_35_Figure_2.jpeg)

Multiple reflections and standing waves will moderate these values

In small primaryfed reflectors multiple reflections between feed and reflector can be used to increase gain

![](_page_36_Picture_1.jpeg)

![](_page_36_Figure_2.jpeg)

# Resonant reflectors can be very efficient and influence system design strongly

![](_page_37_Picture_1.jpeg)

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#### Cassegrain: Subreflector diffraction and blockage

GO is only valid when reflectors are large. Finite diameter of subreflector cause edge diffraction losses. To keep losses smalld >  $10\lambda/\sin(\psi_0)$ 

![](_page_39_Figure_2.jpeg)

![](_page_40_Figure_0.jpeg)

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### System characteristics: Illustration of other subefficiencies

![](_page_42_Figure_1.jpeg)

# System performance: Figure of merit G/T

![](_page_43_Figure_1.jpeg)

![](_page_43_Figure_2.jpeg)

G/T is measured in dB/K.

# System characteristics: Antenna temperature T

![](_page_44_Figure_1.jpeg)

# System characteristics: Figure of merit G/T

![](_page_45_Figure_1.jpeg)

The most common reference point of G and T are at the input of the LNA. G/T is measured in dB/K.

# G/T optimization of same small reflector with dipole-disk feed

![](_page_46_Figure_1.jpeg)

### Surface tolerances

Finite surface accuracy causes phase errors

 $\phi(\theta, \phi) = (1 + \cos\theta) k \Delta z(\theta, \phi) \approx 2k \Delta z(\theta, \phi)$ 

![](_page_47_Figure_3.jpeg)

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### Combination between choke horn and Eleven antenna

![](_page_49_Figure_1.jpeg)

# Some important points

- BOR1 efficiency is important in non-BOR feed geometries
- In BOR1 feeds the co- and cross-polar radiation patterns determines ALL characteristics (phase center, illumination and spillover efficiency)
- Phase center can be uniquely defined from phase pattern in 45 deg plane and the phase efficiency
- Center blockage can give strong dips in G/T
- Airy's disk type feed excitations are possible in Focal Plane Arrays, but not for single feeds
- Combination feeds (choke-horn and Eleven feed) are feasible for minimum relative frequency separation of 1.6

### Eleven feed

![](_page_51_Picture_1.jpeg)